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基于弹性分析法的大型储罐罐壁应力计算

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摘要:运用短圆柱壳挠曲线微分方程, 基于变形光滑连续条件, 建立用于计算大型储罐罐壁应力的弹性分析力学模型, 推导阶梯厚度壳轴向应力的计算公式, 得到详细的计算过程。采用此方法, 以容积为 $15 \times 10^4 \text{ m}^3$ 大型储罐为算例进行验算, 并对有限单元法数值计算结果和现场实测数据进行对比。结果表明, 该方法能够较精确地计算罐壁应力, 可以为罐壁设计与校核提供参考。

关键词:大型储罐; 弹性分析法; 应力计算; 短圆柱壳

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Stress calculation for shells of large storage oil tank based on elastic analysis method

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Abstract: Based on continuity and smoothness conditions of deflection on shell, the mechanical model was established using elastic analysis method for calculating shell stress distribution of large storage tank. The formulas for calculation of hoop and axial stress on stepped thickness shell were deduced by using deflection differential equation of short cylindrical shell. Taking a tank with volume of $15 \times 10^4 \text{ m}^3$ for example, the detailed calculation process of tank shell was provided based on elastic analysis method. Compared the finite element method calculation results with the testing data, the feasibility and effectiveness of the method were verified. The results show that the method can accurately calculate the stress on tank wall, and can provide a reference for the design and checking of the tank wall.

Key words: large storage tank; elastic analysis method; stress calculation; short cylindrical shell

大型储罐是国家战略石油储备的重要装置, 目前国内的大型储罐容积大多数为 $10 \times 10^4 \text{ m}^3$ 。大型储罐的壁厚设计主要是参考原来的设计标准^[1], 然而随着储罐容量的增大, 罐体面临的工况越来越复杂, 必须根据设计参数更为精确地校核储罐在各种工况下的应力及位移情况。储罐罐壁应力的理论计算都是基于薄壳理论, 然而储罐罐壁是非等厚壳体的组合体, 所受载荷随着液位高度而变化, 使理论分析难度增大。从以往的储罐应力方法^[2,4]来看, 长圆

柱壳法^[3]在每层圈板上只考虑单侧边界效应, 组合圆柱壳法^[4]只在最底层圈板上考虑两端边界效应, 两者本质上都是在应力分析法的基础上对某一特定容积储罐应力进行计算的简化形式。虽然简化了计算过程, 但是其通用性有限, 且都未严格评估下节点剪力对整个罐壁结构的影响。笔者从应力分析法的基本假设和边界条件出发, 利用克雷洛夫函数得到挠曲线方程的另一种等效形式, 使得推导和计算过程更为简便, 同时联合弹性-刚性地基梁耦合法计

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算下节点弯矩和剪力,充分考虑下节点弯矩和剪力对罐壁应力分布的影响,给出完整的理论推导和计算过程,并结合实测数据和有限元计算结果进行对比分析,验证本方法计算大型储罐罐壁应力的可行性。

1 弹性分析法理论计算

按照板壳理论中薄壳的定义,储罐罐壁挠度及应力可按一般壳体理论的假设进行理论计算^[6]。

1.1 单层罐壁圈板挠曲线方程的建立

单层储罐罐壁的力学模型如图1所示。由板壳理论^[7]可知,等厚度圆柱壳受轴对称载荷作用时对称变形的挠曲线微分方程为

$$\frac{d^4 y_i}{dx_i^4} + 4\beta_i^4 y_i = Z_i/D_i, \quad (1)$$

其中

$$D_i = \frac{E\delta_i^3}{[12(1-\nu^2)]}, \quad \beta_i = \sqrt[4]{\frac{E\delta_i}{4R^2 D_i}},$$

$$Z_i = P = \gamma(H_{i-1} - x_i).$$

式中, y_i 为第*i*层罐壁挠度,m; Z_i 为沿*x*轴方向变化的垂直于罐壁的载荷, N/m^2 ; D_i 、 β_i 分别为罐壁板的抗弯刚度和特征系数, $N \cdot m, m^{-1}$; E 为钢板弹性模量,Pa; ν 为泊松比; R 为储罐半径,m; δ_i 为第*i*层罐壁的壁厚,m。

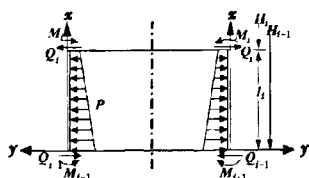


图1 单层罐壁力学模型

Fig.1 Mechanical model of single shell

用克雷洛夫函数表示的轴对称圆柱壳挠曲线方程的通解形式为

$$y_i = C_{1i} V_1(\beta_i x_i) + C_{2i} V_2(\beta_i x_i) + C_{3i} V_3(\beta_i x_i) + C_{4i} V_4(\beta_i x_i) + y^*, \quad (2)$$

其中

$$V_1(\beta_i x_i) = \cosh(\beta_i x_i) \cos(\beta_i x_i),$$

$$V_2(\beta_i x_i) = \frac{1}{2} [\cosh(\beta_i x_i) \sin(\beta_i x_i) + \sinh(\beta_i x_i) \cos(\beta_i x_i)],$$

$$V_3(\beta_i x_i) = \frac{1}{2} \sinh(\beta_i x_i) \sin(\beta_i x_i),$$

$$V_4(\beta_i x_i) = \frac{1}{4} [\cosh(\beta_i x_i) \sin(\beta_i x_i) -$$

$$\sinh(\beta_i x_i) \cos(\beta_i x_i)].$$

式中, C_{1i} 、 C_{2i} 、 C_{3i} 、 C_{4i} 为待定系数。

y^* 是方程(1)的特解, $y^* = \frac{\gamma R^2}{E\delta_i}(H_{i-1} - x_i)$,于是

方程(1)的解可以写成

$$y_i = C_{1i} V_1(\beta_i x_i) + C_{2i} V_2(\beta_i x_i) + C_{3i} V_3(\beta_i x_i) +$$

$$C_{4i} V_4(\beta_i x_i) + \frac{\gamma R^2}{E\delta_i}(H_{i-1} - x_i). \quad (3)$$

$x_i = 0, x_i = h_i$ 处的边界条件为

$$\begin{cases} -D_i \left(\frac{d^2 y_i}{dx_i^2} \right)_{x_i=0} = M_{i-1}, \\ -D_i \left(\frac{d^3 y_i}{dx_i^3} \right)_{x_i=0} = Q_{i-1}, \\ -D_i \left(\frac{d^2 y_i}{dx_i^2} \right)_{x_i=h_i} = M_i, \\ -D_i \left(\frac{d^3 y_i}{dx_i^3} \right)_{x_i=h_i} = Q_i. \end{cases} \quad (4)$$

利用克雷洛夫函数的性质,式(4)可变为

$$\begin{cases} C_{3i} \beta_i^2 = -M_{i-1}/D_i, \\ C_{4i} \beta_i^3 = -Q_{i-1}/D_i, \\ -4C_{1i} \beta_i^2 V_3(\beta_i h_i) - 4C_{2i} \beta_i^2 V_4(\beta_i h_i) + \\ C_{3i} \beta_i^2 V_1(\beta_i h_i) + C_{4i} \beta_i^2 V_2(\beta_i h_i) = -M_i/D_i, \\ -4C_{1i} \beta_i^3 V_2(\beta_i h_i) - 4C_{2i} \beta_i^3 V_3(\beta_i h_i) - \\ 4C_{3i} \beta_i^3 V_4(\beta_i h_i) + C_{4i} \beta_i^3 V_1(\beta_i h_i) = -Q_i/D_i. \end{cases} \quad (5)$$

令 $\alpha_i = \beta_i h_i$,求解得到各项系数为

$$C_{1i} = Z(\alpha_i) [X_i V_3(\alpha_i) - Y_i V_4(\alpha_i)],$$

$$C_{2i} = Z(\alpha_i) [Y_i V_3(\alpha_i) - X_i V_2(\alpha_i)],$$

$$C_{3i} = -\frac{M_{i-1}}{D_i \beta_i^2}, \quad C_{4i} = -\frac{Q_{i-1}}{D_i \beta_i^3}.$$

其中

$$Z(\alpha_i) = \frac{1}{V_3^2(\alpha_i) - V_2(\alpha_i) V_4(\alpha_i)},$$

$$X_i = \frac{M_i}{4D_i \beta_i^2} - \frac{M_{i-1}}{4D_i \beta_i^2} V_1(\alpha_i) - \frac{Q_{i-1}}{4D_i \beta_i^3} V_2(\alpha_i),$$

$$Y_i = \frac{Q_i}{4D_i \beta_i^3} + \frac{M_{i-1}}{D_i \beta_i^2} V_4(\alpha_i) - \frac{Q_{i-1}}{4D_i \beta_i^3} V_1(\alpha_i).$$

于是式(3)的圆柱壳挠度方程可表示为

$$y_i = Z(\alpha_i) [X_i V_3(\alpha_i) - Y_i V_4(\alpha_i)] V_1(\beta_i x_i) + Z(\alpha_i)$$

$$[Y_i V_3(\alpha_i) - X_i V_2(\alpha_i)] V_2(\beta_i x_i) - \frac{M_{i-1}}{D_i \beta_i^2} V_3(\beta_i x_i) -$$

$$\frac{Q_{i-1}}{D_i \beta_i^3} V_4(\beta_i x_i) + \frac{\gamma R^2}{E\delta_i}(H_{i-1} - x_i). \quad (6)$$

1.2 边界条件

对于由 n 层等厚度圆柱壳组合而成的储罐罐壁,相邻圈板交界处的位移和转角相等,因此第 i 层圈板和第 $i-1$ 层圈板交界处的连续条件为

$$y_i \Big|_{x_i=h_i} = y_{i+1} \Big|_{x_{i+1}=0} \quad (7)$$

$$\frac{dy_i}{dx_i} \Big|_{x_i=h_i} = \frac{dy_{i+1}}{dx_{i+1}} \Big|_{x_{i+1}=0} \quad (8)$$

将式(6)代入到式(7)、(8)中,得到

$$\begin{aligned} y_i \Big|_{x_i=h_i} - y_{i+1} \Big|_{x_{i+1}=0} &= Z(\alpha_i) [X_i V_3(\alpha_i) - Y_i V_4(\alpha_i)] \times \\ &V_1(\alpha_i) + Z(\alpha_i) [Y_i V_3(\alpha_i) - X_i V_2(\alpha_i)] V_2(\alpha_i) - \\ &\frac{M_{i-1}}{D_i \beta_i^2} V_3(\beta_i x_i) - \frac{Q_{i-1}}{D_i \beta_i} V_4(\beta_i x_i) - Z(\alpha_{i+1}) [X_{i+1} V_3(\alpha_{i+1}) - \\ &Y_{i+1} V_4(\alpha_{i+1})] + \frac{\gamma R^2}{E \delta_i} (H_{i-1} - H_i) - \frac{\gamma R^2}{E \delta_{i+1}} H_i = 0, \\ \frac{dy_i}{dx_i} \Big|_{x_i=h_i} - \frac{dy_{i+1}}{dx_{i+1}} \Big|_{x_{i+1}=0} &= -4\beta_i Z(\alpha_i) [X_i V_3(\alpha_i) - \\ &Y_i V_4(\alpha_i)] V_4(\alpha_i) + \beta_i Z(\alpha_i) [Y_i V_3(\alpha_i) - X_i V_2(\alpha_i)] \times \\ &V_1(\alpha_i) - \beta_{i+1} Z(\alpha_{i+1}) [Y_{i+1} V_3(\alpha_{i+1}) - X_{i+1} V_2(\alpha_{i+1})] - \\ &\frac{M_{i-1}}{D_i \beta_i} V_2(\alpha_i) - \frac{Q_{i-1}}{D_i \beta_i^2} V_3(\beta_i x_i) - \frac{\gamma R^2}{E \delta_i} + \frac{\gamma R^2}{E \delta_{i+1}} = 0. \end{aligned}$$

上述两等式展开后合并同类项,可以写成如下形式:

$$rm_{i-1} M_{i-1} + rq_{i-1} Q_{i-1} + rm_i M_i + rq_i Q_i + rm_{i+1} M_{i+1} + rq_{i+1} Q_{i+1} = PM_i \quad (9)$$

$$sm_{i-1} M_{i-1} + sq_{i-1} Q_{i-1} + sm_i M_i + sq_i Q_i + sm_{i+1} M_{i+1} + sq_{i+1} Q_{i+1} = PQ_i \quad (10)$$

其中

$$\begin{aligned} rm_{i-1} &= -\frac{Z(\alpha_i)}{4D_i \beta_i^2} V_1^2(\alpha_i) V_3(\alpha_i) - \frac{Z(\alpha_i)}{D_i \beta_i^2} V_4^2(\alpha_i) V_1(\alpha_i) + \\ &\frac{Z(\alpha_i)}{D_i \beta_i^2} V_2(\alpha_i) V_3(\alpha_i) V_4(\alpha_i) + \frac{Z(\alpha_i)}{4D_i \beta_i^2} V_2^2(\alpha_i) V_1(\alpha_i) - \\ &\frac{V_3(\alpha_i)}{D_i \beta_i^2}, \\ rq_{i-1} &= -\frac{Z(\alpha_i)}{2D_i \beta_i^3} V_1(\alpha_i) V_2(\alpha_i) V_3(\alpha_i) + \\ &\frac{Z(\alpha_i)}{4D_i \beta_i^3} V_1^2(\alpha_i) V_4(\alpha_i) + \frac{Z(\alpha_i)}{4D_i \beta_i^3} V_2^2(\alpha_i) - \frac{V_4(\alpha_i)}{D_i \beta_i^3}, \\ rm_i &= \frac{Z(\alpha_i)}{4D_i \beta_i^2} V_1(\alpha_i) V_3(\alpha_i) - \frac{Z(\alpha_i)}{4D_i \beta_i^2} V_2^2(\alpha_i) + \\ &\frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^2} V_1(\alpha_{i+1}) V_3(\alpha_{i+1}) + \frac{Z(\alpha_{i+1})}{D_{i+1} \beta_{i+1}^2} V_4^2(\alpha_{i+1}), \end{aligned}$$

$$\begin{aligned} rq_i &= -\frac{Z(\alpha_i)}{4D_i \beta_i^3} V_1(\alpha_i) V_4(\alpha_i) + \frac{Z(\alpha_i)}{4D_i \beta_i^3} V_2(\alpha_i) V_3(\alpha_i) + \\ &\frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^3} V_2(\alpha_{i+1}) V_3(\alpha_{i+1}) - \frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^3} V_1(\alpha_{i+1}) V_4(\alpha_{i+1}), \\ rm_{i+1} &= -\frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^2} V_3(\alpha_{i+1}), \\ rq_{i+1} &= \frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^3} V_4(\alpha_{i+1}), \\ PM_i &= \frac{\gamma R^2}{E \delta_i} (H_{i-1} - H_i) - \frac{\gamma R^2}{E \delta_{i+1}} H_i, \\ sm_{i-1} &= \frac{Z(\alpha_i)}{D_i \beta_i} V_1(\alpha_i) V_3(\alpha_i) V_4(\alpha_i) + \frac{4Z(\alpha_i)}{D_i \beta_i} V_4^3(\alpha_i) + \\ &\frac{Z(\alpha_i)}{D_i \beta_i} V_1(\alpha_i) V_3(\alpha_i) V_4(\alpha_i) + \frac{Z(\alpha_i)}{4D_i \beta_i} V_1^2(\alpha_i) V_2(\alpha_i) - \\ &\frac{V_2(\alpha_i)}{D_i \beta_i}, \\ sq_{i-1} &= \frac{Z(\alpha_i)}{D_i \beta_i^2} V_2(\alpha_i) V_3(\alpha_i) V_4(\alpha_i) - \frac{Z(\alpha_i)}{D_i \beta_i^2} V_4^2(\alpha_i) \times \\ &V_1(\alpha_i) - \frac{Z(\alpha_i)}{4D_i \beta_i^2} V_1^2(\alpha_i) V_3(\alpha_i) + \frac{Z(\alpha_i)}{4D_i \beta_i^2} V_2^2(\alpha_i) V_1(\alpha_i) - \\ &\frac{V_3(\alpha_i)}{D_i \beta_i^2}, \\ sm_i &= -\frac{Z(\alpha_i)}{D_i \beta_i} V_3(\alpha_i) V_4(\alpha_i) - \frac{Z(\alpha_i)}{4D_i \beta_i} V_1(\alpha_i) V_2(\alpha_i) - \\ &\frac{Z(\alpha_{i+1})}{D_{i+1} \beta_{i+1}} V_4(\alpha_{i+1}) V_3(\alpha_{i+1}) - \frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}} V_1(\alpha_{i+1}) V_2(\alpha_{i+1}), \\ sq_i &= \frac{Z(\alpha_i)}{D_i \beta_i^2} V_4^2(\alpha_i) + \frac{Z(\alpha_i)}{4D_i \beta_i^2} V_1(\alpha_i) V_3(\alpha_i) + \\ &\frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^2} V_1(\alpha_{i+1}) V_3(\alpha_{i+1}) - \frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^2} V_2^2(\alpha_{i+1}), \\ sm_{i+1} &= \frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}} V_2(\alpha_{i+1}), \\ sq_{i+1} &= -\frac{Z(\alpha_{i+1})}{4D_{i+1} \beta_{i+1}^2} V_3(\alpha_{i+1}), \\ PQ_i &= -\frac{\gamma R^2}{E \delta_i} - \frac{\gamma R^2}{E \delta_{i+1}}. \end{aligned}$$

1.3 弯矩和剪力的求解

对于 n 层罐壁,由 $n-1$ 个交界面处的连续条件可以得到由式(9)、(10)确定的 $2(n-1)$ 个等式, $2(n-1)$ 个未知数为 $M_1, Q_1, \dots, M_{n-1}, Q_{n-1}, M_0, Q_0$ 为下节点弯矩和剪力,可由底板应力计算方法^[11]求解, $M_n = Q_n = 0$ 。联立各式得到下列方程组:

得到的储罐下节点应力分布如图 3 (最大应力为 512 MPa,最小应力为 0.106 MPa)所示。

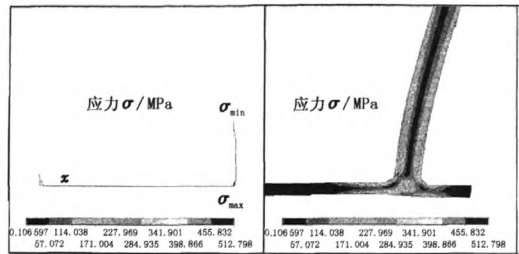


图 3 储罐及其下节点处应力分布云图
Fig.3 Contour of stress distribution on tank and bottom-shell junction

2.4 结果讨论

为了验证弹性分析法计算罐壁应力的合理性,联合实测数据^[12]和有限元数值模拟结果进行对比分析,图 4 为罐壁外侧应力分布(h 为罐壁测点距离罐底上表面的高度)。从图 4 中可以看出,罐壁外侧最大环向应力出现在第 1 层和第 2 层交界处附近。理论计算曲线和实测数据走向基本一致,而且与有限元计算数据非常吻合(罐壁外侧轴向应力也是如此)。表 2 为外壁 6 个测点环向应力理论计算值与实测值对比。从表 2 中可以看出,第 1 层和第 2 层圈板交界处附近 6 个测量点的实测值与本文理论计算值的最大相对误差仅为 5.364%,可以满足工程设计精度要求,且相对误差比组合圆柱壳方法低。

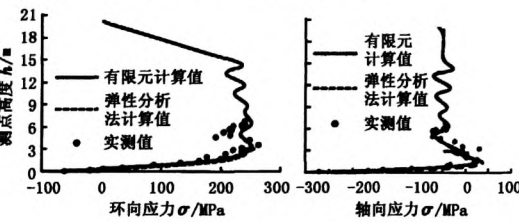


图 4 外壁环向应力及轴向应力分布
Fig.4 Hoop stress and axial stress on outside surface

表 2 外壁 6 个测点环向应力理论计算值与实测值对比
Table 2 Comparison of theoretical values and measuring values of hoop stress at 6 points on outer-shell

测点高度 h/m	环向应力 σ_θ/MPa			相对误差/%	
	实测值	组合圆柱壳法	弹性分析法	组合圆柱壳法	弹性分析法
2.950	251.3	249.9	249.02	-0.557	-0.907
3.006	246.5	251.5	249.63	2.028	1.270
3.097	237.5	251.8	250.24	6.021	5.364
3.102	243.7	251.8	248.89	3.324	2.130
3.390	264.6	250.9	256.32	-5.178	-3.129
4.082	249.9	241.8	242.79	-3.241	-2.845

3 结 论

- (1)引入克雷洛夫函数使得弹性分析法的推导和计算过程变得更为简单,计算结果也较精确,可以将弹性分析法应用于大型储罐设计计算中。
- (2)弹性分析法充分考虑下节点处边缘力系对 各层壁板的应力分布影响,可以较精确地计算大型储罐罐壁应力分布,精度较组合圆柱壳法有所提高。随着储罐的大型化,下节点处应力更为复杂,而下节点处又是整个油罐应力最为复杂的区域,计算时应考虑下节点剪力和弯矩。
- (3)可以利用弹性分析法对变点法或定点法设计的大型储罐罐壁进行应力校核,并根据应力分布情况对罐壁板尺寸进行修正,以达到经济合理设计大型储罐的目的。

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